

Yang-Mills origin of gravitational symmetries

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A. Anastasiou

M.J. Duff, L. Borsten, M. Hughes and S. Nagy

Theoretical Physics Group
Imperial College London

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Outline

Introduce Gravity as the square of Yang-Mills
String Theory origin

Towards a correspondance

Multiplets, Global symmetries, Amplitudes and Solutions

Local symmetries

Squaring

Linear Supergravity and Linear Yang-Mills

The dictionary

Comments and further work

$D = 10$ Type II from SYM

The massless open-string states consist of a vector and a spinor: (A_μ, χ) . Using the $SO(8)$ tensor products:

$$\mathbf{8}_v \otimes \mathbf{8}_v = \mathbf{35}_v \oplus \mathbf{28} \oplus \mathbf{1} \quad , \quad \mathbf{8}_v \otimes \mathbf{8}_s = \mathbf{56}_c \oplus \mathbf{8}_c$$

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 (A_\mu, \chi^+) \otimes (A_\nu, \chi^-) &= (g_{\mu\nu} \oplus B_{\mu\nu} \oplus \phi) + (\psi_\mu^- \oplus \chi^+ \oplus \psi_\mu^+ \oplus \chi^-) \\
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 &\quad + (D_{\mu\nu\rho\sigma}^+ \oplus C_{\mu\nu} \oplus \phi)
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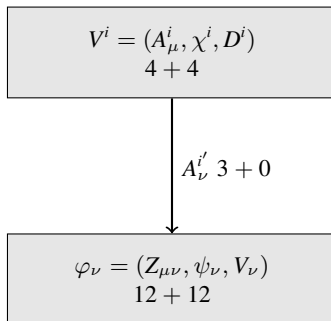
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- **Amplitudes:** Calculating gravity scattering amplitudes using the double-copy formalism. ([Bern, Carrasco, Johansson 0805.3993](#))
- **Solutions:** Schwarzschild is the double copy of the point Coulomb solution. ([Monteiro, O'Connell, White 1410.0239](#))



$(\mathcal{N}_L = 1) \otimes (\mathcal{N}_R = 0)$ gives new-minimal $\mathcal{N} = 1$

Off-shell multiplets



Gravity transformations

By writing the metric as:

$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}$$

the diffeos:

$$x'^{\mu} = x^{\mu} + \kappa \xi^{\mu}$$

give the "gauge" transformation:

$$\delta g_{\mu\nu} = \delta h_{\mu\nu} = \kappa \nabla_{\mu} \xi_{\nu} + \kappa \nabla_{\nu} \xi_{\mu} = \kappa (\partial_{\mu} \xi_{\nu} + \partial_{\nu} \xi_{\mu}) + \mathcal{O}(\kappa^2)$$

and thus to linear order:

$$\delta h_{\mu\nu} = \partial_{\mu} \xi_{\nu} + \partial_{\nu} \xi_{\mu}$$

Yang-Mills transformations

By writing the gauge parameter as:

$$\vartheta^i(x) = \theta^i + g\sigma^i(x)$$

the Yang-Mills transformation becomes:

$$\delta A_\mu^i = \partial_\mu \vartheta^i - g f_{jk}^i \vartheta^j A_\mu^k = g(\partial_\mu \sigma^i - f_{jk}^i \theta^j A_\mu^k) + \mathcal{O}(g^2)$$

and thus to linear order:

$$\delta A_\mu^i = \partial_\mu \sigma^i - f_{jk}^i \theta^j A_\mu^k$$



The goal

The goal is to reproduce the gravity gauge transformations:

$$\begin{aligned}\delta Z_{\mu\nu} &= \partial_\mu \alpha_\nu + \partial_\nu \beta_\mu \quad \text{where } Z_{\mu\nu} \equiv h_{\mu\nu} + B_{\mu\nu} \\ \delta \psi_\nu &= \partial_\nu \eta \\ \delta V_\nu &= \partial_\nu \Lambda\end{aligned}\tag{1}$$

from the Yang-Mills gauge transformations:

$$\begin{aligned}\delta A_\mu^i &= \partial_\mu \sigma^i - f_{jk}^i \theta^j A_\mu^k \\ \delta \chi^i &= -f_{jk}^i \theta^j \chi^k \\ \delta D^i &= -f_{jk}^i \theta^j D^k \\ \delta A_\nu^{i'} &= \partial_\nu \sigma^{i'} - f_{j'k'}^{i'} \theta^{j'} A_\nu^{k'}\end{aligned}\tag{2}$$



The field dictionary

The field dictionary is:

$$\begin{aligned}Z_{\mu\nu} &\equiv A_{\mu}^i \star \Phi_{ii'} \star A_{\nu}^{i'} \\ \psi_{\nu} &\equiv \chi^i \star \Phi_{ii'} \star A_{\nu}^{i'} \\ V_{\nu} &\equiv D^i \star \Phi_{ii'} \star A_{\nu}^{i'}\end{aligned}\tag{3}$$

where the spectator field transforms in the bi-adjoint.

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where the spectator field transforms in the bi-adjoint. **The convolution is crucial in the derivation as it does not obey the Leibniz rule:**

$$\begin{aligned} (a \star b)(x) &= \int a(y) b(x-y) d^4y \\ \partial_\mu (a \star b) &= \partial_\mu a \star b = a \star \partial_\mu b \end{aligned} \tag{4}$$



The transformations

By varying the Yang-Mills parts, the gravity transformations become:

$$\begin{aligned}\delta Z_{\mu\nu} &= \partial_\mu(\sigma^i \star \Phi_{ii'} \star A_{\nu}^{i'}) + \partial_\nu(A_{\mu}^i \star \Phi_{ii'} \star \sigma^{i'}) \\ \delta\psi_{\nu} &= \partial_\nu(\chi^i \star \Phi_{ii'} \star \sigma^{i'}) \\ \delta V_{\nu} &= \partial_\nu(D^i \star \Phi_{ii'} \star \sigma^{i'})\end{aligned}\tag{5}$$



The parameter dictionary

Thus the parameter dictionary is:

$$\begin{aligned}\alpha_\nu &\equiv \sigma^i \star \Phi_{ii'} \star A_\nu^{i'} \\ \beta_\mu &\equiv A_\mu^i \star \Phi_{ii'} \star \sigma^{i'} \\ \eta &\equiv \chi^i \star \Phi_{ii'} \star \sigma^{i'} \\ \Lambda &\equiv D^i \star \Phi_{ii'} \star \sigma^{i'}\end{aligned}\tag{6}$$

Comments

What else works:

- Poincare transformations work trivially.
- Supersymmetry transformations follow in a nice way.

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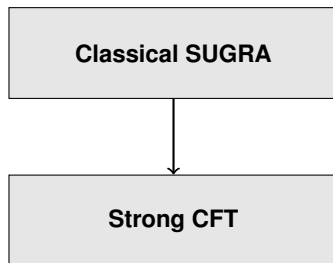
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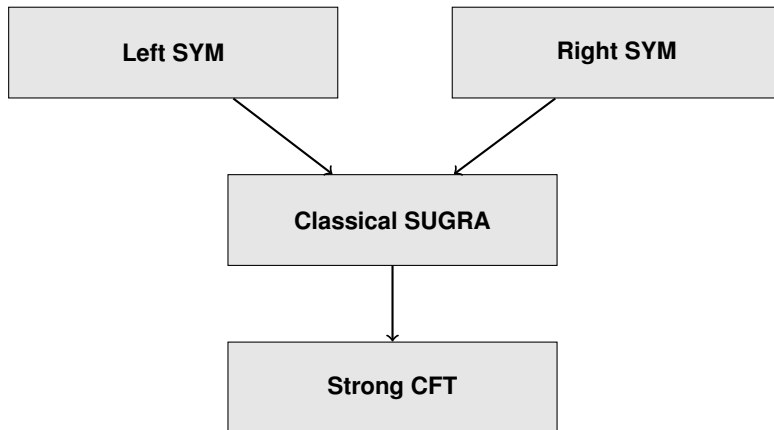
Two major improvements:

- Extend to $\mathcal{N}_R = 1$
- Extend to higher orders where non-linear effects start to appear.

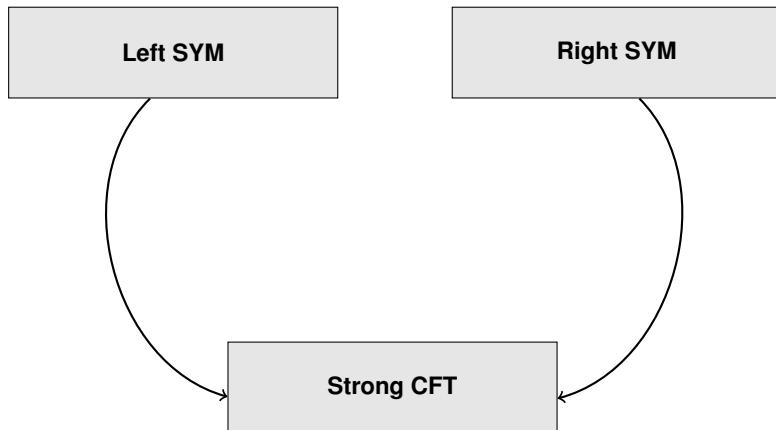
Future work



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THANK YOU